

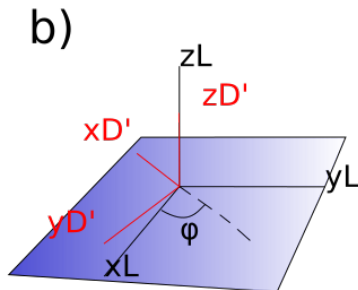
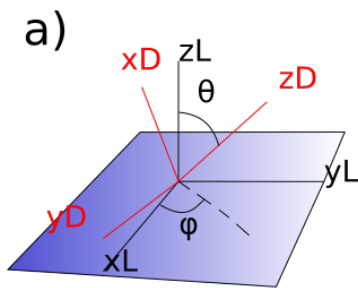
OOP-ESEEM for CISS of photo-generated radical pairs

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Euler angles

- χ axis \equiv dipolar axis \equiv zD
- initial condition: $|\uparrow\downarrow(\mathbf{n})\rangle$
- change of coordinate system: from dipolar frame (D) to laboratory frame (L)



$$R = R_2 \circ R_1,$$
$$R_1 = R_{yD}(\theta), R_2 = R_{zD'}(\pi - \varphi)$$

Euler angles

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$$\begin{aligned} R &= \begin{bmatrix} -\cos(\phi) & \sin(\phi) & 0 \\ -\sin(\phi) & -\cos(\phi) & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \cos(\theta) & 0 & -\sin(\theta) \\ 0 & 1 & 0 \\ \sin(\theta) & 0 & \cos(\theta) \end{bmatrix} \\ &= \begin{bmatrix} -\cos(\theta)\cos(\phi) & \sin(\phi) & \sin(\theta)\cos(\phi) \\ -\cos(\theta)\sin(\phi) & -\cos(\phi) & \sin(\theta)\sin(\phi) \\ \sin(\theta) & 0 & \cos(\theta) \end{bmatrix} \end{aligned}$$

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$$\rho|_D = \rho_{\uparrow\downarrow} = |\uparrow\downarrow\rangle \langle\uparrow\downarrow| = \frac{1}{2} \left(\frac{1}{2}\mathbb{1} - S_{zz} + S_{z,A} - S_{z,B} \right)$$

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$$R_1 = e^{[i(S_{y,A} + S_{y,B})\theta]}, R_2 = e^{[i(S_{z,A} + S_{z,B})(\pi - \phi)]}$$

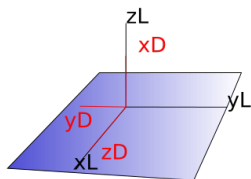
$$\rho|_L = R_2^{-1} R_1^{-1} \rho|_D R_1 R_2$$

Euler angles

$$\begin{aligned}\rho|_L = & -\frac{1}{2}\cos^2(\theta)S_{zz} - \frac{1}{2}\sin^2(\theta)\cos^2(\phi)S_{xx} - \frac{1}{2}\sin^2(\theta)\sin^2(\phi)S_{yy} \\ & + \frac{1}{4}\sin^2(\theta)\sin(2\phi)(S_{xy} + S_{yx}) \\ & + \frac{1}{4}\sin(2\theta)\cos(\phi)(S_{xz} + S_{zx}) - \frac{1}{4}\sin(2\theta)\sin(\phi)(S_{yz} + S_{zy}) \\ & + \frac{1}{2}\cos(\theta)(S_{z,A} - S_{z,B}) \\ & - \frac{1}{2}\sin(\theta)\cos(\phi)(S_{x,A} - S_{x,B}) + \frac{1}{2}\sin(\theta)\sin(\phi)(S_{y,A} - S_{y,B})\end{aligned}$$

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If $\theta = \pi/2$ and $\phi = \pi$:

$$\rho|_L = [-S_{xx} + S_{x,A} - S_{x,B}]/2$$

Hamiltonian and pulse sequence

$$\mathcal{H} = \Omega_A S_{z,A} + \Omega_B S_{z,B} + (d - J) S_{zz} - \frac{d + 2J}{2} (S_{xx} + S_{yy})$$

Hamiltonian and pulse sequence

$$\mathcal{H} = \Omega_A S_{z,A} + \Omega_B S_{z,B} + (d - J) S_{zz} - \frac{d + 2J}{2} (S_{xx} + S_{yy})$$

- Transform to the eigenbasis (therefore \mathcal{H} becomes diagonal):

$$U_\xi = e^{[i(S_{xy} - S_{yx})\xi/2]}$$

$$\mathcal{H} = \Omega_S (S_{z,A} + S_{z,B}) + q' (S_{z,A} - S_{z,B}) + (d - J) S_{zz}$$

with $\xi = \arctan [(d + 2J)/(\Omega_A - \Omega_B)]$, $\Omega_S = (\Omega_A + \Omega_B)/2$,
 $q' = [(\Omega_A - \Omega_B) \cos(\xi) - (d + 2J) \sin(\xi)]/2$.

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In the rotating frame:

$$\mathcal{H} = q' (S_{z,A} - S_{z,B}) + (d - J) S_{zz}$$

Hamiltonian and pulse sequence

$$\mathcal{H} = \Omega_A S_{z,A} + \Omega_B S_{z,B} + (d - J) S_{zz} - \frac{d + 2J}{2} (S_{xx} + S_{yy})$$

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In the rotating frame:

$$\mathcal{H} = q' (S_{z,A} - S_{z,B}) + (d - J) S_{zz}$$

The pulse sequence consists in: lightflash-T- β_x - τ - π_x - τ -echo.

Evolution for isolated operators: part 1

Part 1: lightflash and free evolution for time T .

1. Transform the operator to the eigenbasis of \mathcal{H}
2. Propagation: $\exp[-i\mathcal{H}T]\rho\exp[i\mathcal{H}T]$
3. Fast oscillating approx.: drop all the term \propto sinusoidal of q'

$$\mathcal{H} = \underline{q'(S_{z,A} - S_{z,B})} + (d - J)S_{zz}$$

$$\begin{aligned}\rho|_L = & -\frac{1}{2}\cos^2(\theta)S_{zz} - \frac{1}{2}\sin^2(\theta)\cos^2(\phi)S_{xx} - \frac{1}{2}\sin^2(\theta)\sin^2(\phi)S_{yy} \\ & + \frac{1}{4}\sin^2(\theta)\sin(2\phi)(S_{xy} + S_{yx}) \\ & + \frac{1}{4}\sin(2\theta)\cos(\phi)(\underline{S_{xz} + S_{zx}}) - \frac{1}{4}\sin(2\theta)\sin(\phi)(\underline{S_{yz} + S_{zy}}) \\ & + \frac{1}{2}\cos(\theta)(S_{z,A} - S_{z,B}) \\ & - \frac{1}{2}\sin(\theta)\cos(\phi)(\underline{S_{x,A} - S_{x,B}}) + \frac{1}{2}\sin(\theta)\sin(\phi)(\underline{S_{y,A} - S_{y,B}})\end{aligned}$$

Evolution for isolated operators: part 1

After part 1 of the pulse sequence:

Initial operator

After part 1

$$S_{zz} \rightarrow S_{zz}$$

$$S_{xx} + S_{yy} \rightarrow \sin(\xi)(S_{z,A} - S_{z,B})$$

$$S_{z,A} - S_{z,B} \rightarrow \cos(\xi)(S_{z,A} - S_{z,B})$$

$$S_{xx} - S_{yy} \rightarrow S_{xx} - S_{yy}$$

$$S_{xy} + S_{yx} \rightarrow S_{xy} + S_{yx}$$

$$S_{xz} + S_{zx} \rightarrow 0$$

$$S_{yz} + S_{zy} \rightarrow 0$$

$$S_{x,A} - S_{x,B} \rightarrow 0$$

$$S_{y,A} - S_{y,B} \rightarrow 0$$

Evolution for isolated operators: part 2

Part 2: β_x pulse.

1. Transform back to cartesian basis (inverse of transformation to the eigenbasis of \mathcal{H})
2. Pulse along x -axis with flipping angle β
3. Transform to the eigenbasis of \mathcal{H}

Evolution for isolated operators: part 2

Part 2: β_x pulse.

1. Transform back to cartesian basis (inverse of transformation to the eigenbasis of \mathcal{H})
 2. Pulse along x-axis with flipping angle β
 3. Transform to the eigenbasis of \mathcal{H}
- The initial operators S_{zz} , $S_{xx} + S_{yy}$, $S_{z,A} - S_{z,B}$, $S_{xx} - S_{yy}$:
get terms proportional to: S_{zz} , S_{xx} , S_{yy} , $S_{z,A} - S_{z,B}$
 $S_{yz} + S_{zy}$, $S_{y,A} - S_{y,B}$
 - The initial operator $S_{xy} + S_{yx}$ gets terms
proportional to $S_{xy} + S_{yx}$, $S_{xz} + S_{zx}$, $S_{x,A} - S_{x,B}$

Evolution for isolated operators: part 3

Part 3: free evolution for time τ , π_x -pulse, free evolution for τ .

1. Free evolution (propagation: $\exp[-i\mathcal{H}T]\rho\exp[i\mathcal{H}T]$)
2. Transform to cartesian basis
3. π_x -pulse
4. Transform to eigenbasis of \mathcal{H}
5. Free evolution (propagation: $\exp[-i\mathcal{H}T]\rho\exp[i\mathcal{H}T]$)
6. Fast oscillating approx.: drop all the term \propto sinusoidal of q'
7. Transform to cartesian basis

Evolution for isolated operators: part 3

Part 3: free evolution for time τ , π_x -pulse, free evolution for τ .

1. Free evolution (propagation: $\exp[-i\mathcal{H}T]\rho\exp[i\mathcal{H}T]$)
 2. Transform to cartesian basis
 3. π_x -pulse
 4. Transform to eigenbasis of \mathcal{H}
 5. Free evolution (propagation: $\exp[-i\mathcal{H}T]\rho\exp[i\mathcal{H}T]$)
 6. Fast oscillating approx.: drop all the term \propto sinusoidal of q'
 7. Transform to cartesian basis
 - This part of the sequence does not change the quantum order of the operators
 - To get the signal at the end we do: $\text{Tr}[\rho(S_{x,A} + S_{x,B})]$ and $\text{Tr}[\rho(S_{y,A} + S_{y,B})]$
- We evolve each pathway separately (only single quantum operators)
- We look at the terms proportional to $(S_{x,A} + S_{x,B})$ and $(S_{y,A} + S_{y,B})$ after the evolution

Evolution for isolated operators: part 3

All the pathways that appear at the end of part 2:

<u>S_{zz}</u>	<u>S_{xx}</u>	<u>S_{yy}</u>
<u>$S_{z,A} - S_{z,B}$</u>	$S_{yz} + S_{zy}$	$S_{y,A} - S_{y,B}$
<u>$S_{xy} + S_{yx}$</u>	$S_{xz} + S_{zx}$	$S_{x,A} - S_{x,B}$

- The pathways $S_{yz} + S_{zy}$ and $S_{y,A} - S_{y,B}$ get one term that is proportional to $S_{x,A} + S_{x,B} \rightarrow$ out-of-phase EPR signal
- The pathways $S_{xz} + S_{zx}$ and $S_{x,A} - S_{x,B}$ get one term that is proportional to $S_{y,A} + S_{y,B} \rightarrow$ in-phase EPR signal

Evolution for isolated operators: results

Initial operator	EPR signal
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S_{zz}	\rightarrow out-of-phase
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$S_{xx} + S_{yy}$	\rightarrow out-of-phase
-------------------	----------------------------

$S_{z,A} - S_{z,B}$	\rightarrow out-of-phase
---------------------	----------------------------

$S_{xx} - S_{yy}$	\rightarrow out-of-phase
-------------------	----------------------------

$S_{xy} + S_{yx}$	\rightarrow in-phase
-------------------	------------------------

$S_{xz} + S_{zx}$	$\rightarrow 0$
-------------------	-----------------

$S_{yz} + S_{zy}$	$\rightarrow 0$
-------------------	-----------------

$S_{x,A} - S_{x,B}$	$\rightarrow 0$
---------------------	-----------------

$S_{y,A} - S_{y,B}$	$\rightarrow 0$
---------------------	-----------------

$$\rho_{\uparrow\uparrow/\downarrow\downarrow} = [S_{zz} \pm (S_{z,A} + S_{z,B})]/2 \rightarrow \text{out-of-phase}$$

$$\rho_{T_0/S} = [-S_{zz} \pm (S_{xx} + S_{yy})]/2 \rightarrow \text{out-of-phase}$$

$$\rho_{\uparrow\downarrow/\downarrow\uparrow} = [-S_{zz} \pm (S_{z,A} - S_{z,B})]/2 \rightarrow \text{out-of-phase}$$

Evolution for isolated operators: results

$$E_x(S_{zz}) = \cos^2(\xi) \sin[2(d - J)\tau] \sin(2\beta),$$

$$E_x(S_{xx} + S_{yy}) = \cos^2(\xi) \sin^2(\xi) \sin[2(d - J)\tau][2 \sin(\beta) - \sin(2\beta)],$$

$$E_x(S_{z,A} - S_{z,B}) = \cos^3(\xi) \sin(\xi) \sin[2(d - J)\tau][2 \sin(\beta) - \sin(2\beta)],$$

$$E_x(S_{xx} - S_{yy}) = \sin[2(d - J)\tau] \cos^2(\xi) \sin(2\beta),$$

$$E_y(S_{xy} + S_{yx}) = -2 \cos^2(\xi) \sin(\beta) \sin[2(d - J)\tau].$$

Evolution for isolated operators: results

$$\begin{aligned}\rho|_L = & -\frac{1}{2}\cos^2(\theta)S_{zz} - \frac{1}{2}\sin^2(\theta)\cos^2(\phi)S_{xx} - \frac{1}{2}\sin^2(\theta)\sin^2(\phi)S_{yy} \\ & + \frac{1}{4}\sin^2(\theta)\sin(2\phi)(\underline{S_{xy} + S_{yx}}) \\ & + \frac{1}{4}\sin(2\theta)\cos(\phi)(S_{xz} + S_{zx}) - \frac{1}{4}\sin(2\theta)\sin(\phi)(S_{yz} + S_{zy}) \\ & + \frac{1}{2}\cos(\theta)(S_{z,A} - S_{z,B}) \\ & - \frac{1}{2}\sin(\theta)\cos(\phi)(S_{x,A} - S_{x,B}) + \frac{1}{2}\sin(\theta)\sin(\phi)(S_{y,A} - S_{y,B})\end{aligned}$$

In-phase signal, maximum at
 $\theta = \pi/2$, $\phi = \pi/4$

